

Berman density balance, using bromoform, as $4.60 \pm 0.01 \text{ g.cm}^{-3}$. With the above values of axial lengths and the angle β ($V=1871 \text{ \AA}^3$), the molecular weight ($M=648$ for Fe_7S_8) and the observed density, the number of formula units per unit cell is calculated to be 8 ($Z=7.998$), and the calculated density is 4.602 g.cm^{-3} .

The rotation and Weissenberg photographs and the 15° oscillation photographs about [010] and [001] of a crystal fragment, detached from hexagonal platelets (off an ore sample from level-16), showed superstructure characteristics having a hexagonal supercell $a=6.88$, $c=22.90 \pm 0.02 \text{ \AA}$; the high angle reflexions did not show any separation characteristic of the monoclinic system. The conditions limiting possible reflexions are: $hh2hl$ only with $l=2n$; hhl only with $l=2n$; and $hkil$ with no conditions except $l=2n$; indicating the space group as one of $P6/mcc (D_{6h}^2)$, $P6cc (C_{6v}^2)$. The presence of a number of weak reflexions in the patterns could be explained as being due to the monoclinic phase (see above) which was a minor intergrowth with the hexagonal phase in the crystal fragment. The density for the same crystals was determined as $4.58 \pm 0.01 \text{ g.cm}^{-3}$ and the number of formula units per unit cell was calculated to be 4 ($Z=3.998$) from the above values of axial lengths ($V=939 \text{ \AA}^3$), the observed density and the molecular weight [$M=648$ for Fe_7S_8 although the

composition might have a slight excess of iron (Desborough & Carpenter, 1965)]; the calculated density is 4.582 g.cm^{-3} .

The monoclinic and hexagonal pyrrhotites (superstructure types) were found to be associated mainly with chalcopyrite and pyrite in the ores of different levels (Mukherjee, 1968); the monoclinic phase was predominant in the pyrrhotite-rich ores containing pyrite as a minor constituent, and the hexagonal phase was predominant in the pyrrhotite-poor ores containing chalcopyrite and pyrite as major constituents.

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The Structure of Copper Ammonium Sulfate Hexahydrate from Neutron-Diffraction Data*

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The structure parameters of the Tutton salt $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ have been refined by the method of least squares from three-dimensional neutron-diffraction data. The e.s.d.'s of the atomic coordinates are from 0.0010 to 0.0023 \AA for the N, S, O atoms and 0.0025 to 0.0053 \AA for the H atoms. The value derived for f_s , the coherent scattering amplitude of sulfur, is 0.283(4), significantly different from the value 0.31 in Bacon's table; values derived for f_o and f_N are also slightly different from the tabulated values. Although the heavy-atom coordinates differ significantly in a purely statistical sense from those of the three-dimensional X-ray analysis of Montgomery & Lingafelter, the two descriptions of the heavy-atom structure agree quite well from the point of view of molecular geometry. The hydrogen-bonding pattern from the X-ray study is confirmed.

Introduction

In the last half decade there has been much interest shown, in a number of laboratories, in the crystal structures of Tutton's salts, a well-known series of

isomorphous double sulfates with the general formula $\text{XY}_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$. In an accurate redetermination by X-ray analysis of the structure of $\text{Mg}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$, Margulis & Templeton (1962) were able to locate the hydrogen atoms approximately and to deduce the pattern of hydrogen bonding, after first revising the coordinates of the oxygen atoms from incorrect values reported from the original X-ray determination (Hofmann, 1931). At about the same time and independently, Chidambaram & Rao (1963) con-

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cluded from a proton magnetic resonance (p.m.r.) study of $\text{ZnK}_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ and $\text{MgK}_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ that the oxygen coordinates of Hofmann required revision to allow the assignment of a rational hydrogen-bonding pattern in these salts. There was very satisfying agreement (Templeton, Zalkin & Margulis, 1963) between the orientations of the p - p vectors as determined in the p.m.r. study and in the structure analysis of Margulis & Templeton. The crystal structure of $\text{Mg}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ was also determined independently by Montgomery & Lingafelter (1964*b*), who reported slightly different coordinates for the hydrogen atoms. Now X-ray crystal-structure analyses have been reported for 9 other Tutton's salts, and the complete list of salts analyzed is as follows:

$\text{Mg}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Margulis & Templeton (1962); Montgomery & Lingafelter (1964 <i>b</i>)
$\text{Zn}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Montgomery & Lingafelter (1964 <i>a</i>)
$\text{Ni}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Grimes, Kay & Webb (1963); Montgomery & Lingafelter (1964 <i>b</i>)
$\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Webb, Kay & Grimes (1965); Montgomery & Lingafelter (1966 <i>a</i>)
$\text{Cd}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Montgomery & Lingafelter (1966 <i>b</i>)
$\text{Mn}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	Montgomery, Chastain & Lingafelter (1966)
$\text{V}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	} Montgomery, Chastain, Natt, Witkowska & Lingafelter (1967) Kannan & Viswamitra (1965)
$\text{Fe}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	
$\text{Co}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	
$\text{MgK}_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$	

The results of all these studies are in good general agreement, with slight variations from one salt to another in the structure of the ion $X(\text{H}_2\text{O})_6^{2+}$. It is rather interesting (see Kannan & Viswamitra, 1965) that the tetragonal axis of the g -tensor for each of various Tutton's salts as determined by electron spin resonance (e.s.r.) studies matches one or the other of the vectors $X \rightarrow \text{O}$ in the new structure for Tutton's salts but shows no correlation with the Hofmann structure. This fact went unnoticed by the e.s.r. workers for a decade.

We undertook the neutron-diffraction analysis of the structure of $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ to determine precisely the coordinates of the hydrogen atoms in order to examine in detail the structural parameters of the water molecules, the ammonium ion, and the hydrogen bonds. Soon after we started our study, the results of a two-dimensional X-ray analysis of this salt were published by Webb, Kay & Grimes (1965), and we also

came to know of the three-dimensional analysis that was later published by Montgomery & Lingafelter (1966*a*). Consequently, the comparison of the heavy-atom parameters from our neutron analysis with those from X-ray work became an additional aim of our study.

Experimental

Crystal specimens of appropriate size were grown by slow evaporation from aqueous solution. A crystal in its natural habit weighing 49.9 mg was carefully examined by X-ray precession photography and found to be a single crystal. Approximate cell parameters were obtained from the precession photographs, and the space group symmetry $P2_1/a$ was verified. For the measurements of intensity, the crystal was mounted on a goniometer head with the b axis parallel to the spindle axis.

The recording of the intensity data on the Oak Ridge automatic neutron diffractometer (Busing, Smith, Peterson & Levy, 1964) and the preliminary data processing, by which the raw data were converted to a set of values of observed structure-factor squares and their standard errors, followed the procedures that were used in a recent refinement of the structure of K_2NbF_7 (Brown & Walker, 1966). The same general procedures (Brown & Levy, 1964*b*) have been used recently in a number of other neutron analyses of crystal structures in this laboratory (chloral hydrate, Brown & Levy, 1964*a*; sucrose, Brown & Levy, 1963; α -D-glucose, Brown & Levy, 1965; potassium hydrogen chloromaleate, Ellison & Levy, 1965).

A total of 2250 observations were made in the range $2\theta \leq 111^\circ$ (the limit of the instrument) at a neutron wavelength of 1.078 Å. These included regularly scheduled observations of reference reflections and a few duplicate observations. The number of independent observations was 2032. The few negative values of F_o^2 were replaced by zero values.

For calculation of corrections for absorption (Busing & Levy, 1957; Wehe, Busing & Levy, 1962) the coordinates of the 19 corners of the crystal were measured with a toolmaker's microscope, and the equations of the 12 face planes were calculated. The value 2.2 cm^{-1} for the linear absorption coefficient was determined experimentally. The absorption correction factors were in the range 1.52 to 1.80.

An empirical correction term $(0.03F_o^2)^2$ was added to the statistical variance of each observation F_o^2 to obtain the variance $\sigma^2(F_o^2)$ to be used in computing the weight of the observation in least-squares refinement. The extra term makes allowance for instability in the instrument, deficiencies of the structure model, and so forth (see Peterson & Levy, 1957*a*). The F_o^2 values for 1627 of the 2032 independent reflections were greater than their corrected $\sigma(F_o^2)$ values.

In the course of the collection of the intensity data, the unit-cell parameters were refined, together with three parameters defining the orientation of the crystal

on the diffractometer, by the method of least-squares from the observations of the 2θ positions of the centroids of the diffraction peaks (Busing & Levy, 1967). After our structure refinement was complete, we were provided (Montgomery & Lingafelter, 1966c) with more precise parameters as follows: * $a=9.2105$ (14), $b=12.3795$ (17), $c=6.3016$ (13) Å, $\beta=106.112$ (18)°. These parameters and standard errors were obtained from a least-squares fit to the 2θ values of 17 reflections determined on a diffractometer with Mo $K\alpha$ radiation ($\lambda=0.71069$ Å).

Location of hydrogen atoms and refinement

At the time our data collection was finished, the results of the two-dimensional X-ray analysis of the structure by Webb, Kay & Grimes (1965) had become available. It was clear from these results† that the pattern of hydrogen bonding in the copper ammonium Tutton salt is the typical one that already had been established by the work on the corresponding magnesium, nickel, and zinc salts. Since in this pattern every hydrogen atom of the water molecules and of the ammonium ion is involved in a hydrogen bond to a sulfate oxygen, it was easy to calculate approximate coordinates for the hydrogen atoms. Thus, the structure of Webb *et al.*

* Here and elsewhere in this paper the numbers in parentheses, which correspond to the least significant digits of the adjacent parameters, are the standard errors of the parameters.

† We ignored the conclusion of Webb *et al.* that there is no hydrogen bond from $H_2O(1)$ to $O(4)$, or, in our system of numbering, from $O(8)$ to $O(4')$. Webb *et al.* concluded that the distance 2.54 Å which they calculated for $O(8)\cdots O(4')$ is too short for a hydrogen bond and that there is repulsion between $O(8)$ and $O(4')$. According to them the resultant strain is probably the reason that the copper salt is more liable to dehydration than the corresponding nickel salt. Actually the distance $O(8)\cdots O(4')$ was incorrectly calculated, as were most of the other distances and angles given in their paper; the correct distance corresponding to their coordinates and cell parameters is 2.64 Å. The much more precise X-ray analysis of Montgomery & Lingafelter gave 2.725 (10) Å, a normal value for an $O-H\cdots O$ hydrogen bond. Our neutron analysis has subsequently yielded the distance 2.707(2) Å and has unambiguously shown that there is a hydrogen bond $O(8)\cdots O(4')$. The erroneous conclusion of Webb *et al.* emphasizes that one should be cautious in making inferences about the existence of a hydrogen bond when the hydrogen atom has not been located, and especially cautious when the structural parameters have not been carefully refined.

provided a complete set of starting coordinates for least-squares refinement. Of course, starting coordinates could have been taken equally well from the structure of the magnesium, zinc, or nickel salt.

The neutron coherent scattering amplitudes originally used (Bacon, 1962) were as follows (all in units of 10^{-12} cm): Cu, 0.79; S, 0.31; O, 0.577; N, 0.94; H, -0.378. The structural parameters were refined to convergence (final parameter shift was less than one tenth of the standard error for every parameter) in seven cycles with these scattering factors held fixed. In the last few cycles, anisotropic thermal parameters were optimized. The factor S scaling the $|F_o|$ values to the $|F_c|$ values was also adjusted. In the full-matrix least-squares refinement (Program XFLS, no. 389 in *World List of Crystallographic Computer Programs*, 1966) the quantity minimized was $\sum w(F_o^2 - S^2 F_c^2)^2$, where w is the weight of a particular observation F_o , calculated as $1/\sigma^2(F_o^2)$, and the summation is taken over all symmetrically independent reflections.

Early in the refinement procedure it was clear that many of the F_o^2 values corresponding to observations of high intensity were low because of the effects of extinction, and eventually the 153 observations of highest intensity were omitted from the refinement. Another 81 observations were omitted because their phase signs were indeterminate; these were all reflections with very small F_o^2 and F_c^2 values.

The value of the scale factor and the values of the usual goodness-of-fit measures at this stage are shown in Table 1. The discrepancy indices $R(F^k)$ are defined by the equation $R(F^k) \equiv \sum ||F_o^k| - S^k |F_c^k|| / \sum |F_o^k|$, and the index $R_w(F^2)$ is defined by $R_w(F^2) \equiv [\sum w(F_o^2 - S^2 F_c^2)^2 / \sum w F_o^4]^{1/2}$. The σ_1 is the standard deviation of an observation of unit weight, defined by the equation $\sigma_1 = \sum w(F_o^2 - S^2 F_c^2)^2 / (n - p)$, where w is the weight of an observation F_o^2 , n is the total number of observations, and p is the number of parameters fitted to the data. Unit value is expected for σ_1 at convergence when the observational errors are randomly distributed and correctly estimated and the model is correct.

In two additional cycles, the values of the scattering factors of copper, sulfur, oxygen, and nitrogen were refined to convergence, along with the scale, positional, and thermal parameters. The scattering factor of hydrogen was held at -0.378×10^{-12} cm. The values of the scattering factors derived and their standard errors

Table 1. Scale factors and measures of goodness of fit at convergence in refinement of the structure of $Cu(NH_4)_2(SO_4)_2 \cdot 6H_2O$

See text for definitions.

	Scale factor S	$R(F)$	$R(F^2)$	$R_w(F^2)$	σ_1
Scattering factors not refined*	1.032	0.0865	0.0841	0.0917	1.175
Scattering factors refined, final parameters*	1.022	0.0854	0.0823	0.0896	1.149
Final parameters, data corrected for extinction†	1.022	0.0548	0.0581	0.0765	1.286

* Measures of fit tabulated include only contributions from reflections given non-zero weights in least-squares refinement.

† Measures of fit tabulated include contributions for all reflections observed.

Table 2. Observed neutron structure factors (corrected for absorption but not for extinction) and calculated neutron structure factors for $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$.

For each reflection, identified by the indices H and K of a subcell and the running index L , the values $|F_0|$ and F_0 are given, in units of 10^{-14} cm. The standard error $\sigma(F_0)$, calculated as $\sigma(F_0^2)/2|F_0|$, is given (same units) under the heading SG, except that for each reflection marked W , for which $|F_0|^2 < \sigma(F_0^2)$, the standard error $\sigma(F_0^2)$ is given, in units of 10^{-26} cm. The reflections marked X and U were excluded from the final refinement (see text).

L	OBS	CAL	SG	L	OBS	CAL	SG	L	OBS	CAL	SG	L	OBS	CAL	SG	L	OBS	CAL	SG	
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	
1	220	223	4	1	310	309	4	1	23	200	10	1	280	283	6	1	300	304	7	
2	128	135	5	2	180	181	5	2	109	109	10	2	109	109	10	2	109	109	10	
3	358	360	5	3	250	250	5	3	88	88	10	3	88	88	10	3	88	88	10	
4	46	79	21	4	40	41	21	4	140	140	21	4	140	140	21	4	140	140	21	
5	45	74	21	5	103	103	21	5	111	111	21	5	111	111	21	5	111	111	21	
6	60	40	24	6	127	127	24	6	111	111	24	6	111	111	24	6	111	111	24	
7	70	40	23	7	61	60	23	7	226	226	23	7	226	226	23	7	226	226	23	
8	282	269	7	8	551	576	7	8	307	300	7	8	307	300	7	8	307	300	7	
9	266	261	7	9	347	354	7	9	154	143	7	9	154	143	7	9	154	143	7	
10	10	21	22	10	180	182	22	10	285	295	22	10	285	295	22	10	285	295	22	
11	372	372	6	11	192	183	6	11	168	168	6	11	168	168	6	11	168	168	6	
12	157	151	5	12	87	85	5	12	185	195	5	12	185	195	5	12	185	195	5	
13	149	155	6	13	68	35	22	13	142	128	10	13	142	128	10	13	142	128	10	
14	176	190	7	14	96	93	10	14	101	107	10	14	101	107	10	14	101	107	10	
15	185	182	7	15	35	26	29	15	301	286	7	15	301	286	7	15	301	286	7	
16	0	16	23	16	0	22	32	16	294	113	14	16	294	113	14	16	294	113	14	
17	239	230	6	17	37	34	20	17	221	225	9	17	221	225	9	17	221	225	9	
18	62	90	27	18	13	13	22	18	230	228	8	18	230	228	8	18	230	228	8	
19	200	186	9	19	110	83	12	19	163	160	10	19	163	160	10	19	163	160	10	
20	0	2	22	20	0	21	26	20	166	166	22	20	166	166	22	20	166	166	22	
21	0	2	22	21	0	22	32	21	101	101	22	21	101	101	22	21	101	101	22	
22	228	245	3	22	126	126	30	22	126	126	30	22	126	126	30	22	126	126	30	
23	124	256	4	23	36	47	34	23	112	112	34	23	112	112	34	23	112	112	34	
24	7	21	25	24	0	21	25	24	129	116	4	24	129	116	4	24	129	116	4	
25	99	113	9	25	146	142	10	25	194	202	9	25	194	202	9	25	194	202	9	
26	44	47	48	26	83	74	11	26	125	153	14	26	125	153	14	26	125	153	14	
27	116	114	11	27	358	363	7	27	106	106	11	27	106	106	11	27	106	106	11	
28	146	163	9	28	179	171	10	28	282	281	6	28	282	281	6	28	282	281	6	
29	7	200	186	29	146	152	13	29	426	434	7	29	426	434	7	29	426	434	7	
30	0	13	14	30	44	36	22	30	75	48	19	30	75	48	19	30	75	48	19	
31	0	13	14	31	64	80	23	31	111	122	12	31	111	122	12	31	111	122	12	
32	1	216	214	32	1	216	214	32	1	216	214	32	1	216	214	32	1	216	214	32
33	192	189	8	33	50	49	43	33	222	222	43	33	222	222	43	33	222	222	43	
34	303	303	8	34	55	82	37	34	37	19	33	34	37	19	33	34	37	19	33	
35	106	109	17	35	106	109	17	35	106	109	17	35	106	109	17	35	106	109	17	
36	4	78	73	36	4	78	73	36	4	78	73	36	4	78	73	36	4	78	73	36
37	95	101	13	37	95	101	13	37	95	101	13	37	95	101	13	37	95	101	13	
38	6	240	236	38	6	240	236	38	6	240	236	38	6	240	236	38	6	240	236	38
39	7	352	344	39	7	352	344	39	7	352	344	39	7	352	344	39	7	352	344	39
40	169	169	10	40	169	169	10	40	169	169	10	40	169	169	10	40	169	169	10	
41	0	156	158	41	0	156	158	41	0	156	158	41	0	156	158	41	0	156	158	41
42	150	156	12	42	150	156	12	42	150	156	12	42	150	156	12	42	150	156	12	
43	0	146	145	43	0	146	145	43	0	146	145	43	0	146	145	43	0	146	145	43
44	0	146	145	44	0	146	145	44	0	146	145	44	0	146	145	44	0	146	145	44
45	0	146	145	45	0	146	145	45	0	146	145	45	0	146	145	45	0	146	145	45
46	0	146	145	46	0	146	145	46	0	146	145	46	0	146	145	46	0	146	145	46
47	0	146	145	47	0	146	145	47	0	146	145	47	0	146	145	47	0	146	145	47
48	0	146	145	48	0	146	145	48	0	146	145	48	0	146	145	48	0	146	145	48
49	0	146	145	49	0	146	145	49	0	146	145	49	0	146	145	49	0	146	145	49
50	0	146	145	50	0	146	145	50	0	146	145	50	0	146	145	50	0	146	145	50
51	0	146	145	51	0	146	145	51	0	146	145	51	0	146	145	51	0	146	145	51
52	0	146	145	52	0	146	145	52	0	146	145	52	0	146	145	52	0	146	145	52
53	0	146	145	53	0	146	145	53	0	146	145	53	0	146	145	53	0	146	145	53
54	0	146	145	54	0	146	145	54	0	146	145	54	0	146	145	54	0	146	145	54
55	0	146	145	55	0	146	145	55	0	146	145	55	0	146	145	55	0	146	145	55
56	0	146	145	56	0	146	145	56	0	146	145	56	0	146	145	56	0	146	145	56
57	0	146	145	57	0	146	145	57	0	146	145	57	0	146	145	57	0	146	145	57
58	0	146	145	58	0	146	145	58	0	146	145	58	0	146	145	58	0	146	145	58
59	0	146	145	59	0	146	145	59	0	146	145	59	0	146	145	59	0	146	145	59
60	0	146	145	60	0	146	145	60	0	146	145	60	0	146	145	60	0	146	145	60
61	0	146	145	61	0	146	145	61	0	146	145	61	0	146	145	61	0	146	145	61
62	0	146	145	62	0	146	145	62	0	146	145	62	0	146	145	62	0	146	145	62
63	0	146	145	63	0	146	145	63	0	146	145	63	0	146	145	63	0	146	145	63
64	0	146	145	64	0	146	145	64	0	146	145	64	0	146	145	64	0	146	145	64
65	0	146	145	65	0	146	145	65	0	146	145	65	0	146	145	65	0	146	145	65
66	0	146	145	66	0	146	145	66	0	146	145	66	0	146	145	66	0	146	145	66
67	0	146	145	67	0															

Table 3. Final atomic coordinates and thermal parameters (with standard errors in parentheses) of the crystal structure of $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$

The thermal parameters β_{ij} are the coefficients of the index products in the temperature factor:

$$\exp [-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)].$$

The second line associated with each heavy atom contains, for each parameter, the values Δ , defined as the X-ray parameter value minus the neutron value, and $|\Delta|/\sigma$ (see text). Each coordinate Δ is expressed in units $\text{\AA} \times 10^{-3}$; each Δ for the parameters β_{ij} is on the same scale as the β_{ij} .

	Fractional coordinates $\times 10^5$			Thermal parameters $\times 10^5$					
	<i>x</i>	<i>y</i>	<i>z</i>	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Cu(1)	0 (—)	0 (—)	0 (—)	438 (13)	252 (08)	1054 (29)	0 (07)	207 (14)	31 (10)
				156 9.4	92 10.2	796 13.8	-6 0.8	73 3.7	19 1.5
S(2)	41069 (24)	13898 (19)	74517 (36)	435 (26)	299 (16)	1177 (60)	-37 (14)	208 (27)	-72 (21)
	-17 7.0	6 2.1	4 1.6	-32 1.2	-60 3.6	193 2.3	-13 0.9	-14 0.5	28 1.2
O(3)	41576 (18)	23173 (11)	60057 (23)	1184 (19)	376 (08)	1721 (33)	-198 (10)	514 (20)	-3 (13)
	-7 2.0	4 1.1	2 0.4	-34 0.7	-54 2.9	139 0.9	7 0.3	-25 0.4	9 0.3
O(4)	54888 (16)	7580 (14)	78118 (30)	595 (15)	604 (11)	3296 (53)	104 (10)	221 (22)	-215 (19)
	-9 2.7	7 1.8	4 0.8	-27 0.7	-51 2.1	-116 0.7	-13 0.6	0 0.0	-24 0.5
O(5)	28057 (13)	7032 (11)	63349 (21)	605 (13)	390 (08)	1747 (31)	-137 (08)	291 (16)	-122 (12)
	-4 1.5	-2 0.6	-3 0.7	2 0.0	-41 2.2	-157 1.2	-9 0.5	48 1.0	56 1.6
O(6)	39147 (17)	17862 (12)	95686 (21)	1152 (19)	420 (08)	1376 (29)	-47 (09)	437 (19)	-105 (12)
	-9 2.4	-10 2.8	-2 0.4	-62 1.4	-14 0.7	154 1.1	-3 0.2	-50 0.9	-6 0.2
O(7)	17658 (16)	11673 (12)	17726 (23)	789 (16)	541 (10)	1559 (35)	-31 (10)	290 (19)	61 (14)
	-14 4.2	-22 5.9	2 0.4	-34 0.9	5 0.2	121 0.8	24 1.0	47 0.9	91 2.3
O(8)	-16392 (15)	10892 (12)	3043 (23)	687 (15)	385 (09)	1892 (34)	-24 (09)	278 (18)	92 (13)
	-4 1.4	9 2.8	11 2.6	4 0.1	-48 2.5	408 2.9	-34 1.7	-5 0.1	58 1.6
O(9)	-531 (14)	-6534 (10)	28210 (19)	737 (14)	332 (08)	1309 (27)	28 (08)	405 (15)	112 (11)
	0 0.0	15 5.1	-1 0.2	-92 2.7	-11 0.6	451 3.4	30 1.6	11 0.2	-12 0.4
N(10)	13458 (10)	34786 (08)	35953 (14)	861 (12)	448 (07)	1845 (26)	5 (07)	485 (12)	6 (09)
	-10 2.7	-2 0.6	2 0.5	-103 2.3	-48 2.1	415 2.2	22 1.0	-76 1.2	81 2.0
H(11)	6554 (46)	33321 (38)	20966 (60)	1964 (61)	1334 (40)	2694 (99)	-526 (40)	361 (63)	-226 (49)
H(12)	22515 (43)	29916 (31)	40162 (76)	1565 (51)	801 (27)	5929 (167)	365 (30)	1439 (74)	371 (53)
H(13)	7617 (49)	33427 (42)	46879 (68)	2018 (59)	1573 (45)	3549 (108)	681 (44)	1712 (71)	870 (61)
H(14)	16861 (49)	42467 (30)	36184 (66)	2068 (62)	586 (22)	4149 (125)	-105 (30)	119 (68)	140 (41)
H(15)	22103 (31)	9620 (25)	32899 (41)	1220 (34)	717 (20)	1796 (68)	-31 (20)	193 (37)	103 (27)
H(16)	25555 (31)	12868 (24)	10680 (45)	1150 (33)	731 (20)	2523 (72)	-184 (21)	725 (40)	-62 (29)
H(17)	-26718 (26)	9535 (22)	-5970 (43)	794 (28)	600 (18)	2610 (70)	-12 (16)	292 (35)	8 (28)
H(18)	-14202 (28)	18405 (21)	217 (44)	1113 (32)	423 (16)	2690 (72)	-24 (17)	597 (38)	81 (25)
H(19)	-10271 (28)	-5877 (21)	31563 (41)	1039 (31)	596 (17)	2341 (65)	131 (17)	848 (35)	147 (26)
H(20)	2457 (27)	-14082 (19)	30950 (38)	1055 (28)	431 (16)	2138 (59)	99 (17)	523 (32)	155 (23)

omitted in the final refinement because of extinction errors. The reflections marked *U* are the ones omitted because it is considered that their signs are undetermined.

The final atomic parameters and their standard errors appear in Table 3. A drawing (Program *ORTEP*, no. 387 in *World List*, 1966) of the crystal structure viewed in the direction of $-c$ is shown in Fig. 1.

Discussion

Comparison of neutron and X-ray parameters

Because of the possibility of learning something about the effects of chemical bonding on atomic electron distributions (see Dawson, 1964, 1965; Coppens & Coulson, 1967; Coppens, 1967), we consider it desirable to make a detailed comparison of the parameters from our neutron analysis with those from the three-dimensional X-ray analysis of Montgomery & Lingafelter (hereafter abbreviated to M. & L.). We therefore present in Table 3 the differences between the M. & L. heavy-atom coordinates and thermal parameters and

ours. Each difference Δ in the table, X-ray parameter minus neutron parameter, is accompanied by the ratio of its magnitude to its standard error σ , where σ is computed as $[\sigma_{\text{neutron}}^2 + \sigma_{\text{X-ray}}^2]^{1/2}$. Our Table 4 of bond lengths and angles* involving the heavy atoms only includes for comparison the corresponding values from the M. & L. X-ray work. For the hydrogen atoms, the parameters of which M. & L. do not claim to be accurate, we note only that the positions are different by from 0.05 \AA to 0.4 \AA (average difference 0.24 \AA).

From the point of view of molecular geometry only, the descriptions of the hydrated copper ion and the sulfate ion are in reasonable agreement with those of M. & L. (see Table 4). In particular, the copper-oxygen framework of the cation clearly does have approximate orthorhombic symmetry rather than the essentially te-

* All our bond lengths and angles and functions of thermal parameters, shown in Table 4 and subsequent tables, were computed with Program *ORFFE* 11, modified by Dr C. K. Johnson from *ORFFE* (no. 363 in *World List*, 1966). The full least-squares covariance matrix was used in computing standard errors.

Table 4. Bond lengths and angles not involving the hydrogen atoms in the crystal structure of $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$

The corresponding bond lengths and angles from the X-ray work of Montgomery & Lingafelter are shown for comparison. Bond lengths marked with asterisks are 'corrected' values ('riding' model, see text).

Atoms	Bond length (Å) or angle (°)	
	This work	M. & L.†
Cu(1)–O(7)	2.2300 (14)	2.209 (5)
	2.2369*	
Cu(1)–O(8)	2.0725 (14)	2.084 (5)
	2.0871*	
Cu(1)–O(9)	1.9660 (11)	1.960 (5)
	1.9710*	
S(2)–O(3)	1.475 (3)	1.479 (5)
	1.489*	
S(2)–O(4)	1.457 (3)	1.463 (5)
	1.481*	
S(2)–O(5)	1.479 (3)	1.476 (5)
	1.486*	
S(2)–O(6)	1.477 (3)	1.465 (5)
	1.490*	
O(7)–Cu(1)–O(8)	88.90 (05)	88.89 (15)
O(7)–Cu(1)–O(9)	90.65 (05)	90.35 (15)
O(8)–Cu(1)–O(9)	88.99 (05)	88.61 (15)
O(3)–S(2)–O(4)	109.8 (2)	109.4 (2)
O(3)–S(2)–O(5)	108.3 (2)	108.4 (2)
O(3)–S(2)–O(6)	109.4 (2)	110.1 (2)
O(4)–S(2)–O(5)	108.9 (2)	108.5 (2)
O(4)–S(2)–O(6)	110.9 (2)	110.5 (2)
O(5)–S(2)–O(6)	109.7 (2)	110.2 (2)

† The bond lengths and angles in this column are slightly different from those given in the M. & L. (1966a) paper, which were calculated before the revised unit-cell parameters (M. & L., 1966c) were available.

trigonal symmetry reported by Webb *et al.* from the two-dimensional X-ray analysis. From the point of view of physical statistics, however, we conclude that the small differences between the two sets of heavy-atom coordinates are significant. By application of the *R*-factor ratio test (Hamilton, 1965) the hypothesis that the X-ray heavy-atom coordinates fit the neutron data is most decisively rejected* ($R = 1.275$, $R_{27, 1618, 0.005} = 1.015$, probability $\ll 0.005$). This result implies either that the differences are real effects related to bonding and electron distributions or that in one or the other, or both, of the two determinations data have been included which are affected by systematic errors.

Several considerations suggest that the differences are mainly the results of errors. First, we note that for the *x* direction in the crystal all of the non-zero Δ values except one have the same sign. Neither this fact nor the apparently random signs of the differences in bond lengths is explicable on the basis of bonding effects on electron distributions.

We note on careful examination of the fit of the calculated to the observed structure factors in the structure-factor table of M. & L. that for by far the most of the reflections included in the refinement having $|F_o| \geq 40$ the quantity $|F_o| - |F_c|$ is negative, often by as much as 10% of $|F_c|$. This fact might suggest more pervasive effects of extinction than were recognized by M. & L. However, even in the case of the 12 reflections

* Our thermal parameters were used with the X-ray coordinates in the structure-factor calculation required for determining \mathcal{R} .

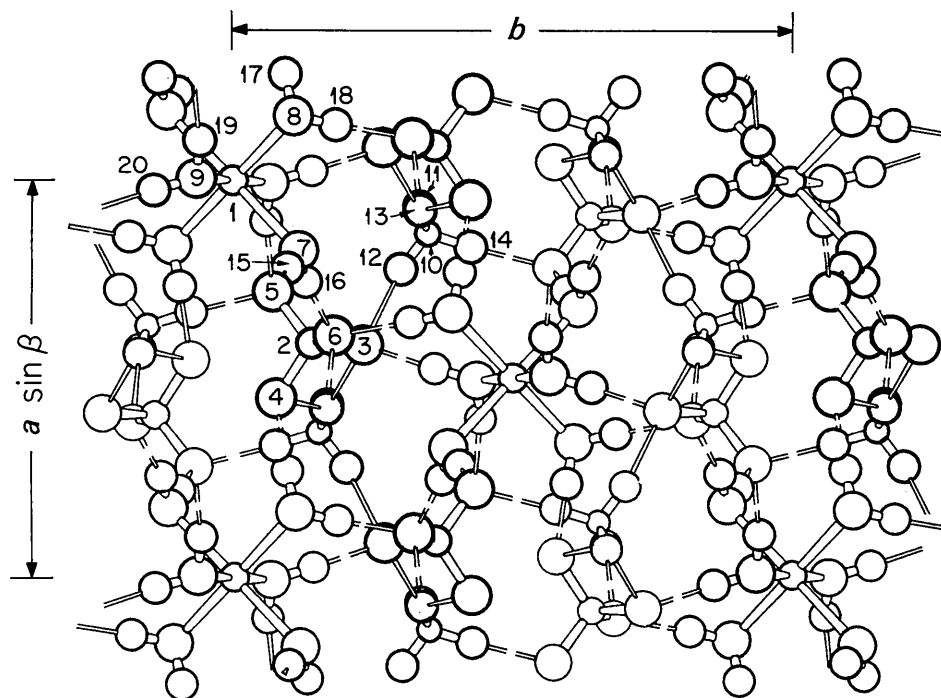


Fig. 1. The crystal structure of $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$, viewed in projection in the direction of $-c$. The atoms of the reference asymmetric unit are numbered as in Table 3. Identification of atoms: 1, Cu; 2, S; 3–9, O; 10, N; 11–20, H.

assumed to be affected by extinction and omitted from the refinement, there is not the expected correlation of the magnitude of $|F_o| - |F_c|$ with the calculated intensity, some of the largest discrepancies being associated with reflections of only moderate intensity. We believe that systematic errors of some other kind may be involved as well and that these other errors in the data may account for the main differences between the X-ray and neutron coordinates. M. & L. do not in fact claim great accuracy in their structure determination.

We cannot assert, of course, that our results are entirely free of the effects of systematic errors. However, most of our standard errors in the coordinates are only about half as large as M. & L.'s and some are about one third as large. Only for the sulfur atom, for which the neutron scattering factor is relatively small and the X-ray scattering factor is moderately large, are the X-ray standard errors smaller (one half to one third of the neutron errors); our errors for this atom are smaller than the X-ray errors for any other atom.

In contrast to the case of copper ammonium sulfate hexahydrate, it has been found in this laboratory (El Saffar & Brown, 1968) that the heavy-atom coordinates from the X-ray determination of structure of $\text{MnCl}_2 \cdot 4\text{H}_2\text{O}$ (counter data, Mo $K\alpha$ radiation, Zalkin, Forrester & Templeton, 1964) fit the neutron data as well as the neutron coordinates do, according to the R -factor ratio test, the probability of significant difference in fit being ~ 0.5 . In the comparison for $\text{MnCl}_2 \cdot 4\text{H}_2\text{O}$, the average and maximum values of $|\Delta|$ for the 21 heavy-atom coordinates are only 0.0021 \AA and 0.0039 \AA , and the average and maximum $|\Delta|/\sigma$ values are 0.8 and 1.9 .

Even without the effects of systematic errors, any attempts at interpreting the thermal parameters of M. & L. would be expected to be difficult because the X-ray data for the layers recorded in Weissenberg photography (five layers hkl with $l=0$ to 4) were not relatively scaled, except approximately, before the least-squares refinement. Therefore, concerning the comparison of neutron and X-ray thermal parameters (Table 4) we only remark that the agreement is surprisingly good, except for the three β_{ii} parameters of copper, for each of which Δ/σ is ~ 10 . These large errors are clearly the result of the fact that the dispersion correction $\Delta f'$ (-2.1 electrons for the Cu atom and Cu $K\alpha$ radiation) was not used in the X-ray analysis. We have found after performing three more cycles of refinement with the dispersion correction and the published X-ray data that the Δ values of the Cu β_{ii} parameters are reduced approximately to the level of the Δ values of the remaining thermal parameters. There was no accompanying significant improvement in the agreement for other parameters.

Thermal motion

Principal-axis transformations were performed on the thermal-parameter ellipsoids. Table 5 gives the

principal root-mean-square vibrational displacements and the orientations of the principal axes in the cell. For the oxygen and hydrogen atoms the orientation of each principal axis is specified with respect to the vector $X \rightarrow A$, where X is the heavier atom to which the given lighter atom A (oxygen or hydrogen) is attached. The root-mean-square displacements of X and A in the direction $X \rightarrow A$ are also given.

The data in Table 5 for O(7), O(8), and O(9) are relevant to consideration of the problem of correcting the Cu–O bond lengths for thermal motion. There is a correlation between the looseness of binding of oxygen to copper as measured by the length of the Cu–O bond and the looseness as measured by the difference between the r.m.s. displacements of the two atoms along the bond direction. The most tightly bound oxygen atom, O(9), is the only one of the three water oxygens for which the principal axis of smallest r.m.s. displacement lies close to the Cu–O bond direction. The orientations of the ellipsoids from the M. & L. X-ray structure are close to those of Table 5. The situation is different from that found in the X-ray study of the magnesium Tutton salt (Margulis & Templeton, 1962) in which the principal axis of smallest displacement of each water oxygen lies close to the Mg–O bond direction. Table 5 includes 'corrected' bond lengths for the Cu–O bonds, calculated as the mean Cu–O separations on the assumption that the oxygen atoms 'ride' the copper atom (Busing & Levy, 1964). The corrected bond-length Cu–O(9) is probably reliable, but the others may not be. Fortunately, the desired corrections are probably small and of little significance in any case.

We judge from the relevant data in the last three columns of Table 5 that the riding model is probably a good approximation for calculating corrected S–O bond lengths. It is interesting that the corrections we calculate are almost identical with the riding corrections for the S–O bonds calculated by Margulis & Templeton (1962) in their X-ray study of the isomorphous magnesium ammonium sulfate hexahydrate. No bond or angle difference between our sulfate ion and that in the magnesium compound is greater than one standard deviation (theirs) and almost every one is considerably less. Our average values of 1.473 \AA and 1.486 \AA for the uncorrected and corrected S–O bonds are nearly identical with theirs.

Water molecules, ammonium ion, and hydrogen bonding

It is convenient to discuss the water molecules and the ammonium ion in relation to the hydrogen bonding. Our results confirm the pattern of hydrogen bonding in $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$ as found by Montgomery & Lingafelter (1966a) and at the same time furnish a much more precise description of the bonding. Table 6 shows the interatomic distances and angles which describe the water molecules, the ammonium ion, and the hydrogen bonding. The hydrogen bonding is nearly symmetrical on the two sides of each water molecule.

Table 5. Description of atomic thermal-parameter ellipsoids of $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$: orientations of the principal axes and root-mean-square displacements (*r.m.s.d.*)

The orientations of the principal axes (p.a.) are specified with respect to the unit vectors of a Cartesian coordinate system and to interatomic vectors $X \rightarrow A$. The vector *i* is along the crystal vector *a*; *j* is along *b*; *k* is along *c**. Atom *X* is the heavier atom to which a given listed atom *A* is attached.

Atom	p.a.	R.m.s.d. (Å)	Angle (°) between p.a. and reference axis				R.m.s.d. (Å) of <i>X</i> and <i>A</i> along $X \rightarrow A$	
			<i>i</i>	<i>j</i>	<i>k</i>	$X \rightarrow A$	<i>X</i>	<i>A</i>
Cu(1)	1	0.130 (2)	15 (13)	79 (10)	99 (10)			
	2	0.136 (2)	75 (13)	132 (10)	46 (10)			
	3	0.144 (2)	91 (06)	44 (09)	46 (09)			
S(2)	1	0.129 (4)	21 (11)	74 (08)	77 (10)			
	2	0.142 (4)	70 (12)	122 (09)	140 (09)			
	3	0.160 (4)	96 (05)	37 (08)	127 (08)			
O(3)	1	0.151 (2)	61 (01)	33 (02)	105 (03)	25 (2)	0.157 (4)	
	2	0.177 (2)	86 (02)	110 (03)	160 (02)	72 (3)	0.162 (2)	
	3	0.229 (2)	29 (01)	116 (01)	77 (01)	107 (1)		
O(4)	1	0.149 (2)	19 (01)	104 (01)	77 (01)	19 (1)	0.143 (4)	
	2	0.209 (2)	96 (01)	151 (02)	118 (02)	72 (1)	0.157 (2)	
	3	0.264 (2)	108 (01)	115 (02)	32 (01)	85 (1)		
O(5)	1	0.138 (2)	38 (02)	56 (02)	75 (02)	167 (2)	0.132 (4)	
	2	0.175 (2)	58 (02)	115 (03)	137 (04)	78 (2)	0.140 (2)	
	3	0.193 (2)	108 (03)	44 (03)	129 (04)	85 (2)		
O(6)	1	0.152 (2)	101 (01)	69 (02)	23 (02)	11 (1)	0.143 (4)	
	2	0.183 (2)	105 (02)	157 (02)	73 (02)	87 (2)	0.154 (2)	
	3	0.215 (2)	19 (02)	101 (02)	75 (01)	101 (1)		
O(7)	1	0.169 (2)	95 (11)	101 (03)	12 (07)	73 (9)	0.139 (2)	
	2	0.176 (2)	167 (05)	101 (03)	98 (11)	140 (6)	0.187 (2)	
	3	0.207 (2)	102 (03)	15 (02)	81 (02)	55 (3)		
O(8)	1	0.162 (2)	34 (10)	56 (10)	90 (07)	105 (9)	0.136 (2)	
	2	0.170 (2)	62 (11)	134 (09)	57 (03)	139 (6)	0.179 (2)	
	3	0.195 (2)	108 (03)	63 (03)	33 (03)	53 (3)		
O(9)	1	0.139 (2)	66 (03)	61 (03)	141 (02)	170 (2)	0.136 (2)	
	2	0.161 (2)	48 (04)	137 (04)	96 (04)	88 (4)	0.140 (2)	
	3	0.179 (2)	52 (04)	61 (04)	51 (02)	80 (2)		
N(10)	1	0.171 (1)	37 (02)	91 (03)	127 (02)			
	2	0.186 (2)	86 (06)	173 (09)	84 (07)			
	3	0.193 (1)	54 (02)	83 (09)	37 (03)			
H(11)	1	0.214 (4)	71 (03)	73 (02)	26 (04)	173 (2)	0.191 (1)	
	2	0.259 (4)	140 (03)	119 (02)	65 (04)	91 (4)	0.216 (4)	
	3	0.351 (5)	124 (02)	34 (02)	94 (02)	83 (1)		
H(12)	1	0.200 (4)	142 (03)	52 (03)	91 (02)	167 (2)	0.185 (1)	
	2	0.262 (4)	125 (03)	137 (03)	68 (02)	85 (3)	0.209 (4)	
	3	0.342 (5)	78 (02)	72 (02)	22 (02)	78 (1)		
H(13)	1	0.189 (4)	46 (02)	88 (02)	136 (02)	171 (1)	0.171 (1)	
	2	0.249 (4)	54 (03)	126 (01)	57 (03)	88 (3)	0.196 (4)	
	3	0.394 (6)	66 (01)	36 (01)	64 (01)	98 (1)		
H(14)	1	0.210 (4)	83 (03)	9 (02)	95 (03)	13 (4)	0.186 (2)	
	2	0.254 (4)	123 (03)	91 (04)	147 (03)	101 (4)	0.214 (4)	
	3	0.327 (5)	34 (03)	99 (02)	123 (03)	84 (2)		
H(15)	1	0.179 (4)	79 (03)	98 (02)	13 (02)	8 (3)	0.170 (2)	
	2	0.225 (3)	147 (08)	123 (08)	85 (03)	85 (3)	0.180 (4)	
	3	0.242 (3)	121 (08)	34 (08)	78 (02)	95 (2)		
H(16)	1	0.185 (3)	32 (03)	71 (03)	114 (04)	11 (3)	0.175 (2)	
	2	0.218 (3)	103 (04)	117 (04)	149 (04)	84 (4)	0.187 (3)	
	3	0.252 (3)	118 (02)	35 (04)	108 (04)	99 (2)		
H(17)	1	0.177 (3)	19 (03)	87 (03)	71 (03)	163 (3)	0.168 (2)	
	2	0.216 (3)	90 (06)	172 (15)	82 (15)	85 (5)	0.182 (3)	
	3	0.225 (3)	109 (03)	82 (15)	20 (07)	107 (3)		
H(18)	1	0.178 (4)	76 (05)	18 (05)	101 (03)	1 (3)	0.167 (2)	
	2	0.206 (3)	160 (06)	73 (05)	79 (06)	90 (5)	0.178 (4)	
	3	0.226 (3)	76 (06)	82 (03)	16 (05)	88 (3)		
H(19)	1	0.163 (3)	33 (02)	96 (03)	122 (03)	160 (2)	0.159 (2)	
	2	0.201 (3)	76 (03)	143 (03)	57 (03)	101 (3)	0.172 (3)	
	3	0.239 (3)	61 (02)	53 (03)	50 (03)	107 (2)		
H(20)	1	0.172 (3)	101 (06)	29 (03)	116 (05)	163 (5)	0.157 (2)	
	2	0.192 (3)	141 (05)	82 (06)	52 (06)	102 (7)	0.175 (4)	
	3	0.217 (3)	53 (05)	62 (03)	49 (04)	101 (3)		

The lone-pair coordination of the water molecules is of type *J* in the classification scheme of Chidambaram, Sequeira & Sikka (1964); that is, one of the lone pairs of electrons on each water molecule is directed approximately toward a polyvalent metal ion and the other lone pair is not specifically directed. The two-lone pair orbitals and the two OH bonds are considered to make approximately regular tetrahedral angles one with another (Coulson, 1957). The O-H bond lengths before correction are all in the range 0.961–0.980 Å, the shortest two being in the water molecule which has the weakest hydrogen bonding. The corrected bond lengths range from 0.982 to 1.003 Å, averaging 0.991 Å. The riding model does not actually apply exactly here, and consequently the corrections are probably over-corrections. The valence angle 105.9 (3)° for the water molecule H(17)–O(8)–H(18) and the angle 105.7 (3)° for molecule H(19)–O(9)–H(20) are close to the value 104.52 (5)° reported for water vapor (Benedict, Gailar

& Plyler, 1956; for earlier values see references in *Tables of, Interatomic Distances* 1958), but the angle 109.3 (3)° for molecule H(15)–O(7)–H(16) is distinctly higher.

The water molecule with the largest valence angle, H(15)–O(7)–H(16), is the one having the largest associated O...O...O angle (118.04°), the weakest (that is, longest) hydrogen bonds, and the longest Cu–O coordination distance. The large O...O...O angle suggests that some part of the increase of the valence angle H–O–H from the vapor value of 104.5° must result from the forces exerted on the hydrogen atoms along the lines H...O. However, most of the deformation of the valence angle must be attributed to some other cause, since Chidambaram (1962) has shown that bending of the O–H...O bonds is energetically much more favorable than distortion of the H–O–H angle. If the O...O...O angle were an important factor, one would think that the water molecule H(19)–O(9)–H(20)

Table 6. Distances and angles which describe the water molecules, the ammonium ion, and the hydrogen bonding in $\text{Cu}(\text{NH}_4)_2(\text{SO}_4)_2 \cdot 6\text{H}_2\text{O}$

The bond lengths *X–H* in the table marked with asterisks are 'corrected' values ('riding' model, see text). The atom number marked with asterisks indicate atoms related to the atoms of the asymmetric unit by symmetry transformations, which are specified in the paper of Montgomery & Lingafelter (1966a).

Hydrogen bond <i>X–H</i> ... <i>O</i>	Distances (Å)			Angles (°)		
	<i>X–H</i>	<i>H</i> ... <i>O</i>	<i>X</i> ... <i>O</i>	<i>X–H</i> ... <i>O</i>	<i>H</i> ... <i>O–S</i>	<i>H–O–Cu</i>
O(7)–H(15)···O(5)	0.964 (3) 0.982*	1.871 (3)	2.826 (2)	170.1 (3)	111.0 (2)	112.1 (2)
O(7)–H(16)···O(6*)	0.963 (3) 0.983*	1.871 (3)	2.821 (2)	168.7 (3)	127.7 (2)	114.1 (2)
O(8)–H(17)···O(4*)	0.978 (3) 0.993*	1.730 (3)	2.707 (2)	178.1 (3)	133.5 (2)	116.8 (2)
O(8)–H(18)···O(6*)	0.978 (3) 0.993*	1.765 (3)	2.743 (2)	177.5 (3)	122.4 (2)	114.2 (2)
O(9)–H(19)···O(5*)	0.981 (3) 0.991*	1.760 (3)	2.732 (2)	170.4 (3)	129.2 (2)	114.5 (2)
O(9)–H(20)···O(3*)	0.976 (3) 1.002*	1.715 (3)	2.683 (2)	171.3 (2)	125.1 (2)	118.5 (2)
N(10)–H(11)···O(6*)	0.998 (4) 1.061*	1.927 (4)	2.899 (2)	164.1 (4)	116.1 (2)	
N(10)–H(12)···O(3)	1.004 (4) 1.061*	2.032 (4)	2.983 (2)	157.3 (4)	122.1 (2)	
O(3*)		2.055 (4)	3.007 (2)	158.6 (5)	95.2 (2)	
N(10)–H(13)···O(4*)	0.999 (4) 1.070*					
N(10)–H(14)···O(5)	1.000 (4) 1.052*	2.335 (5) 1.861 (4)	3.121 (2) 2.860 (2)	134.9 (4) 176.3 (4)	84.6 (2) 112.5 (2)	

H– <i>X</i> –H angles (°)	Distance <i>H</i> ... <i>H</i> (Å)	O... <i>X</i> ...O angles (°)
H(15)–O(7)–H(16)	1.572 (4)	O(5)···O(7)···O(6*) 118.04 (6)
H(17)–O(8)–H(18)	1.561 (4)	O(4*)···O(8)···O(6*) 104.17 (7)
H(19)–O(9)–H(20)	1.560 (3)	O(5*)···O(9)···O(3*) 98.98 (6)
H(11)–N(10)–H(12)	1.678 (6)	O(6*)···N(10)···O(3) 137.87 (5)
H(11)–N(10)–H(13)	1.609 (5)	O(6*)···N(10)···O(3*) 86.67 (5)
		O(6*)···N(10)···O(4*) 117.58 (5)
H(11)–N(10)–H(14)	1.612 (5)	O(6*)···N(10)···O(5*) 105.29 (5)
H(12)–N(10)–H(13)	1.604 (5)	O(3)···N(10)···O(3*) 101.18 (4)
		O(3)···N(10)···O(4*) 95.82 (5)
H(12)–N(10)–H(14)	1.636 (5)	O(3)···N(10)···O(5*) 105.15 (4)
H(13)–N(10)–H(14)	1.658 (6)	O(3*)···N(10)···O(5) 121.77 (4)
		O(4*)···N(10)···O(5) 80.09 (5)

should have a valence angle smaller than it actually has, since the associated angle $O \cdots O \cdots O$ is in this case only 98.98° and the hydrogen bonds involved are stronger than those of $H(15)-O(7)-H(16)$.

The bonding of a water molecule to a copper ion may be expected to affect the structure of the water molecule. Thus, Peterson & Levy (1957*b*) have noted that in cupric chloride dihydrate the planarity of the entire grouping $CuCl_2 \cdot 2H_2O$, including the H atoms, suggests π character in the Cu–O bonds. It also suggests an accompanying increase of *s* character in the O–H bonds, consistent with the observed opening of the H–O–H angle from the vapor value to 108° (Chidambaram, 1963; Chidambaram, Sequeira & Sikka, 1964; see also Pimentel & McClellan, 1960, and Frank, 1958). The lone-pair coordination in $CuCl_2 \cdot 2H_2O$ is of type *D* in the classification scheme of Chidambaram *et al.* (1964); that is, the bisector of the angle between the directions of the lone pairs is directed toward the cupric ion. For coordination of type *J*, as in our case, the effect on the valence angle would be expected to be smaller than for type *D*, consistent with the angles observed for molecules $H(17)-O(8)-H(18)$ and $H(19)-O(9)-H(20)$ but not with the angle of molecule $H(15)-O(7)-H(16)$. Other hydrates in which rather large H–O–H angles have been found in neutron-diffraction studies are $CuSO_4 \cdot 5H_2O$ (Bacon & Curry, 1962; see also the nuclear magnetic resonance (n.m.r.) study of El Saffar, 1966) and $MnCl_2 \cdot 4H_2O$ (El Saffar & Brown, 1968). In the latter hydrate there is a correlation between the coordination types and the H–O–H angles of the four water molecules: two type *D* molecules have angles of $104.3 (0.3)^\circ$ and $106.4 (0.2)^\circ$, and two of type *J* have angles of $112.4 (0.3)^\circ$ and $111.4 (0.3)^\circ$. In $CuSO_4 \cdot 5H_2O$ the angles $114 (1)^\circ$ and $109 (1)^\circ$ of two type *D* molecules are not much different from the angles $111 (1)^\circ$ and $109 (1)^\circ$ of two type *J* molecules; the fifth water molecule, which is not coordinated to the copper ion, has the angle $106 (1)^\circ$. It is clear that more data are needed for understanding the variations in the shape of the water molecule in hydrates. Also needed is a rational procedure for correcting the valence angle for the effects of thermal motion, which must limit our understanding to some degree at present.

In a comparison of n.m.r. and neutron-diffraction results for twelve crystalline hydrates, El Saffar (1966) has included a comparison of the magnitudes (see Table 6) and orientations of the *p*–*p* vectors of the water molecules from our study of $Cu(NH_4)_2(SO_4)_2 \cdot 6H_2O$ with the average values from the two sets of vectors obtained for $K_2Zn(SO_4)_2 \cdot 6H_2O$ and $K_2Mg(SO_4)_2 \cdot 6H_2O$ in the n.m.r. study of Chidambaram & Rao (1963). If allowance is made for the fact that the same Tutton salt was not used in the two techniques, the agreement seems quite satisfactory.

The accuracy of the determination of the N–H bond lengths is impaired by the effects of the considerable thermal motion of the hydrogen atoms of the ammonium ion. The average N–H bond length before cor-

rection, 1.000 \AA , becomes 1.06 \AA after application of the riding model correction to the individual bonds, which we judge should be a good approximation in this case. These values are close to the corresponding values from neutron-diffraction studies reported by Schlemper, Hamilton & Rush (1966): $0.985 (7)$ and $1.064 (7) \text{ \AA}$ for the uncorrected and corrected values in $(NH_4)_2SiF_6$; $1.06 (1) \text{ \AA}$ for the mean corrected value in $(NH_4)_2SO_4$ (uncorrected values not given). It is difficult, unfortunately, to assess the accuracy of bond lengths corrected by use of the riding model when the corrections amount to 0.06 to 0.08 \AA , as in these cases. The close agreement in the N–H bond lengths from $Cu(NH_4)_2(SO_4)_2 \cdot 6H_2O$ and $(NH_4)_2SiF_6$ may be fortuitous, especially in view of the disordered arrangement of H atoms which complicated the diffraction study of the latter compound.

Other values reported previously for the N–H bond length in ammonium salts determined by neutron analysis have been consistently lower than 1.06 \AA : $1.03 (2) \text{ \AA}$ in NH_4Cl and ND_4Cl (Levy & Peterson, 1952); $1.026 (2)$ to $1.037 (2) \text{ \AA}$ in various forms of ND_4Br (Levy & Peterson, 1953*a*); 1.03 \AA in ND_4I (Levy & Peterson, 1953*b*); 1.004 \AA in $NH_4H_2PO_4$ (Tenzer, Frazer & Pepinsky, 1958); $1.03 (3) \text{ \AA}$ in $(NH_4)_2BeF_4$ (Kay, Lajzerowicz, Okaya & Pepinsky, 1959); 1.023 \AA (average) in $(NH_4)_2C_2O_4 \cdot H_2O$ (Padmanabhan, Srikantha & Ali, 1965); 1.02 \AA and 1.044 \AA (corrected), both averages, in HN_3OHCl (Padmanabhan, Smith & Peterson, 1967); 1.03 \AA (average) in $LiN_2H_5SO_4$ (Padmanabhan & Balasubramanian, 1967). Most of these neutron determinations were based on a small number of intensity data*, and in most of them the refinements were not so complete as the present one. In general the effects of thermal motion on bond lengths were not taken into account.† It is not clear, however, whether the bond length 1.06 \AA from $Cu(NH_4)_2(SO_4)_2 \cdot 6H_2O$, $(NH_4)_2SiF_6$, and $(NH_4)_2SO_4$ should be regarded as more reliable than values from the other neutron-diffraction studies. It should be noted here that N–H bond lengths of 1.03 to 1.04 \AA have been reported from n.m.r. work on NH_4Br and NH_4Cl (Gutowsky, Pake & Bersohn, 1954; Bersohn & Gutowsky, 1954).

The hydrogen bonds $N-H \cdots O$ are, on the average, bent more than the bonds $O-H \cdots O$. Atom H(13) is the hydrogen involved in what was described by Montgomery & Lingafelter as a bifurcated hydrogen bond. The distance $H(13) \cdots O(4)$ of 2.332 \AA indicates a rather weak interaction, and the main strength of the bifurcated hydrogen bond lies in the $H(13) \cdots O(3)$ interaction. There appears to be a rough correlation

* The NH_3OHCl determination was the only one based on even partial three-dimensional data.

† In the work by Levy & Peterson on the ammonium halides the effects of thermal motion on bond lengths were allowed for explicitly in the structure analysis; however, the structure parameters were determined only by trial calculations.

between the angles H–N–H, which range from 106.5° to 113.9°, and the corresponding angles O...N...O.

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